

## CONTROL OF QUADRATURE

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The use of the midpoint rule in the treatment of the right-hand side in the Poisson problem, i.e., the approximation

$$\ell_h(v_h) = \sum_{T \in \mathcal{T}_h} |T| f(x_T) v_h(x_T)$$

of the exact functional

$$\ell(v_h) = \int_{\Omega} f v_h \, dx$$

leads to an error contribution that is controlled by the first Strang lemma. Defining the elementwise constant function  $f_h \in L^2(\Omega)$  by requiring that

$$f_h|_T = f(x_T)$$

for all  $T \in \mathcal{T}_h$  and noting that the midpoint rule is exact for affine functions we have that

$$\ell_h(v_h) = \sum_{T \in \mathcal{T}_h} f(x_T) \int_T v_h \, dx = \int_{\Omega} f_h v_h \, dx.$$

This implies that

$$\ell(v_h) - \ell_h(v_h) = \int_{\Omega} (f - f_h) v_h \, dx.$$

With Hölder's inequality we find that

$$\ell(v_h) - \ell_h(v_h) \leq \|f - f_h\|_{L^2(\Omega)} \|v_h\|_{L^2(\Omega)}.$$

Assuming that  $f \in W^{1,\infty}(\Omega)$  or that  $f$  is Lipschitz continuous it follows from the mean value theorem that for  $T \in \mathcal{T}_h$  and  $x \in T$  we have

$$|f(x) - f(x_T)| \leq h_T \|\nabla f\|_{L^\infty(T)}.$$

With this estimate we deduce that

$$\|\ell(v_h) - \ell_h(v_h)\|_{V_h'} \leq h \|\nabla f\|_{L^\infty(\Omega)}.$$

The argument also applies if  $f$  is only elementwise Lipschitz continuous. A quadratic error contribution can be obtained if the midpoint rule is applied with the nodal interpolant  $\mathcal{I}_h f$  of  $f$ , i.e., if we use

$$\ell_h(v_h) = \sum_{T \in \mathcal{T}_h} |T| \mathcal{I}_h f(x_T) v_h(x_T) = \int_{\Omega} f_h v_h \, dx,$$

where  $f_h$  is the elementwise constant function with  $f_h|_T = \mathcal{I}_h f(x_T)$  for every  $T \in \mathcal{T}_h$ . Since  $\mathcal{I}_h f$  is elementwise affine we have that

$$\int_T \mathcal{I}_h f \, dx = |T| \mathcal{I}_h f(x_T) = \int_T f_h \, dx,$$

which implies that

$$\int_T (\mathcal{I}_h f - f_h) v_h \, dx = \int_T (\mathcal{I}_h f - f_h) (v_h - \alpha_T) \, dx$$

for every  $\alpha_T \in \mathbb{R}$ . We let  $\bar{v}_h$  be the elementwise constant function with

$$\bar{v}_h|_T = \frac{1}{|T|} \int_T v_h \, dx$$

for every  $T \in \mathcal{T}_h$ . We thus have that

$$\begin{aligned} \ell(v_h) - \ell_h(v_h) &= \int_{\Omega} (f - \mathcal{I}_h f) v_h \, dx + \int_{\Omega} (\mathcal{I}_h f - f_h) v_h \, dx \\ &= \int_{\Omega} (f - \mathcal{I}_h f) v_h \, dx + \int_{\Omega} (\mathcal{I}_h f - f_h) (v_h - \bar{v}_h) \, dx \\ &\leq \|f - \mathcal{I}_h f\|_{L^2(\Omega)} \|v_h\|_{L^2(\Omega)} + \|\mathcal{I}_h f - f_h\|_{L^2(\Omega)} \|v_h - \bar{v}_h\|_{L^2(\Omega)}. \end{aligned}$$

Using nodal interpolation estimates we find that  $\|f - \mathcal{I}_h f\| \leq ch^2 \|D^2 f\|$ . To control the second term we note that for every  $v \in W^{1,2}(T)$  with vanishing integral we have that

$$\|v\|_{L^2(T)} \leq c_T h_T \|\nabla v\|_{L^2(T)}$$

with a constant  $c_T$  that depends on the geometry of the element  $T$  and which is uniformly bounded if  $\mathcal{T}_h$  is shape-regular. This leads to the estimate

$$\|\mathcal{I}_h f - f_h\|_{L^2(\Omega)} \leq ch \|\nabla \mathcal{I}_h f\|_{L^2(\Omega)}, \quad \|v_h - \bar{v}_h\|_{L^2(\Omega)} \leq ch \|\nabla v_h\|_{L^2(\Omega)},$$

and leads to

$$|\ell(v_h) - \ell_h(v_h)| \leq ch^2 \|D^2 f\|_{L^2(\Omega)} \|v_h\|_{L^2(\Omega)} + ch^2 \|\nabla \mathcal{I}_h f\|_{L^2(\Omega)} \|\nabla v_h\|_{L^2(\Omega)}.$$

Interpolation estimates and the triangle inequality show that

$$\|\nabla \mathcal{I}_h f\|_{L^2(\Omega)} \leq \|\nabla f\|_{L^2(\Omega)} + ch \|D^2 f\|_{L^2(\Omega)}.$$

Combining the estimates we find that

$$\|\ell - \ell_h\|_{V'_h} \leq ch^2 \|f\|_{H^2(\Omega)}.$$

The calculations show that it is useful to choose an elementwise polynomial function  $f_h$  such that the discrete functional

$$\ell_h(v_h) = \int_{\Omega} f_h v_h \, dx$$

can be computed exactly.